



Micromodelling of polyethylene materials

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PRESENTATION OUTLINE

1. **Research objectives**
2. Modelling of crystalline polyethylene
3. Modelling of semicrystalline polyethylene
4. Conclusions

PROBLEM STATEMENT

Motivation

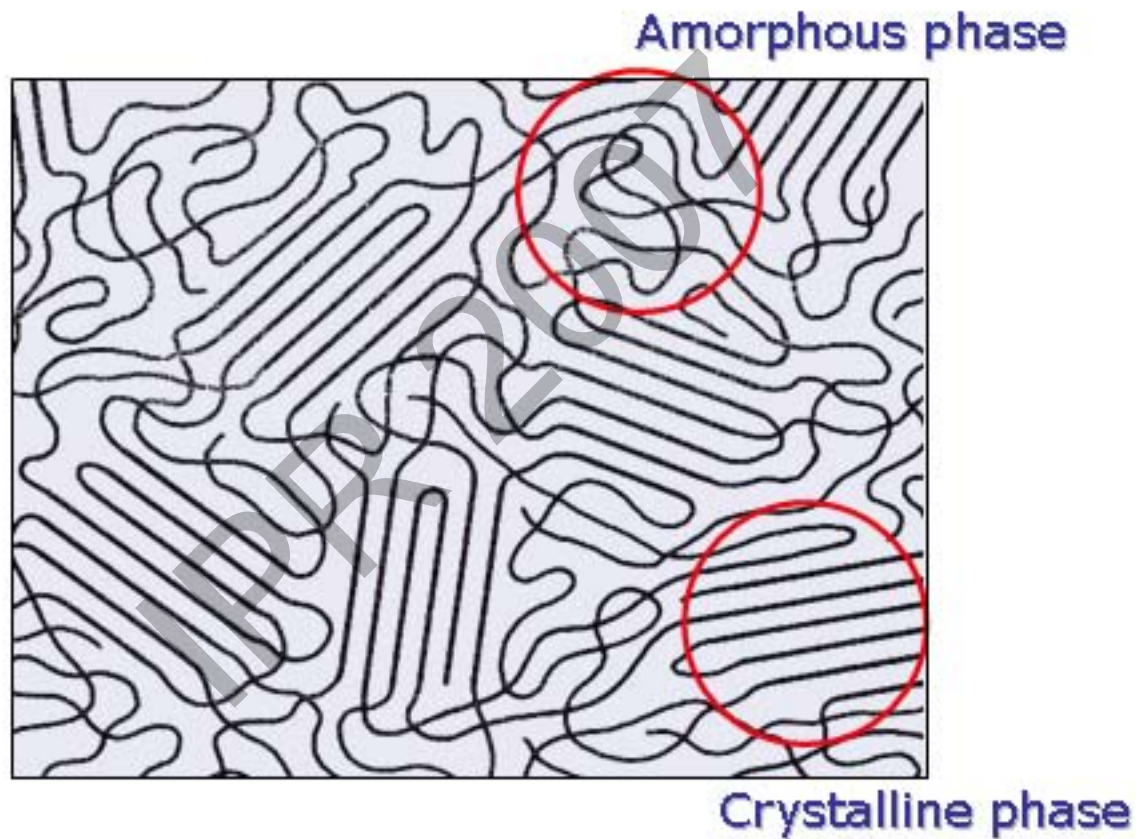
- To understand the interplay between the microstructure and the overall macroscopic behaviour of semicrystalline polyethylene.

Goal

- To understand the deformation process of polyethylene
- To develop constitutive models to predict the mechanical behaviour of polyethylene considering the damage processes occurring at large deformations.

POLYETHYLENE PHASES

Polyethylene structure

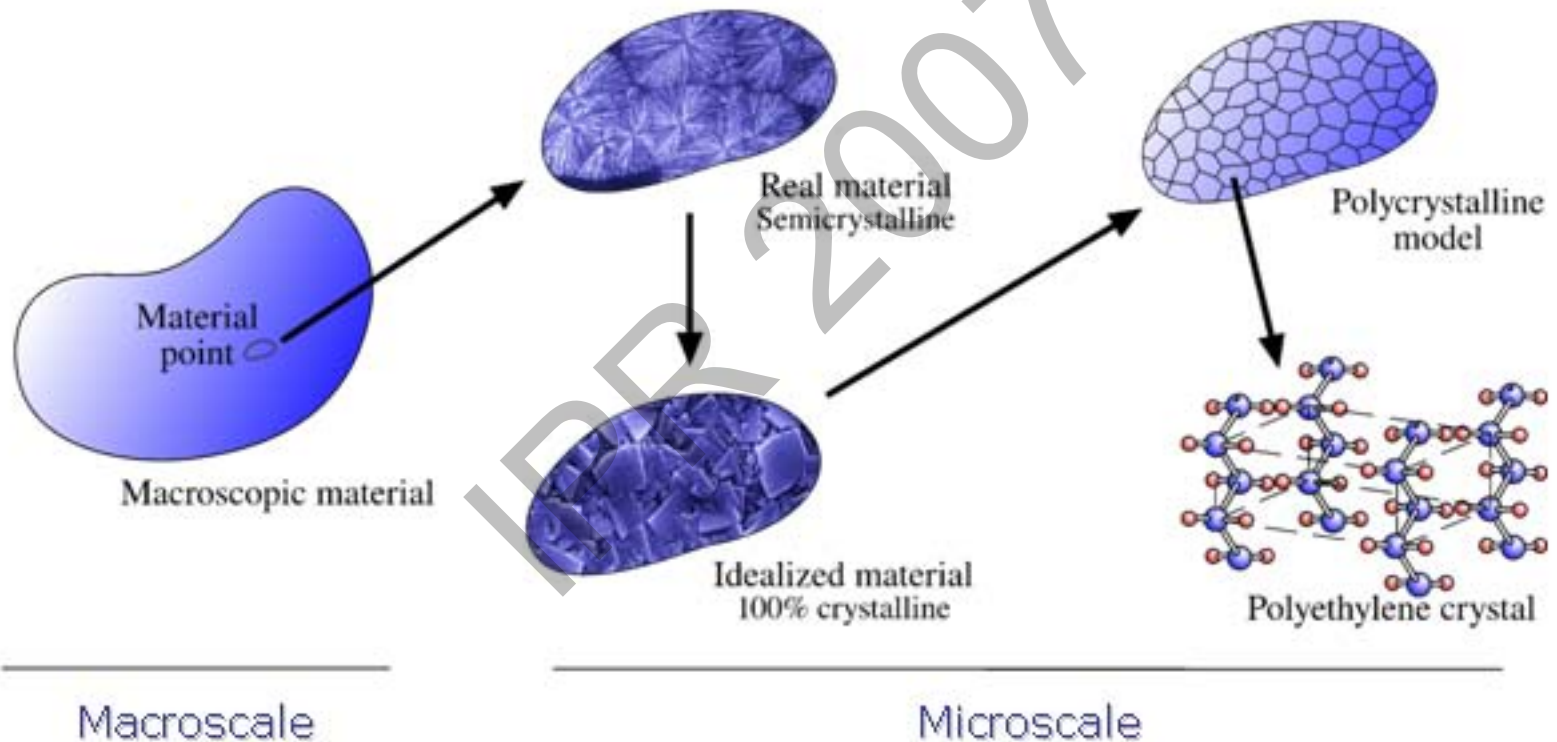


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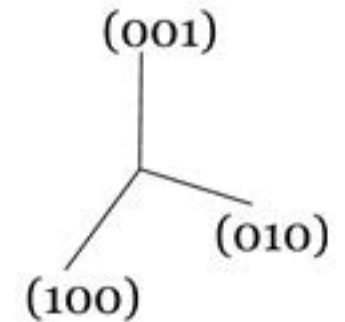
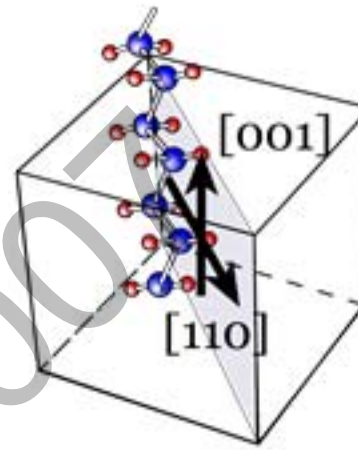
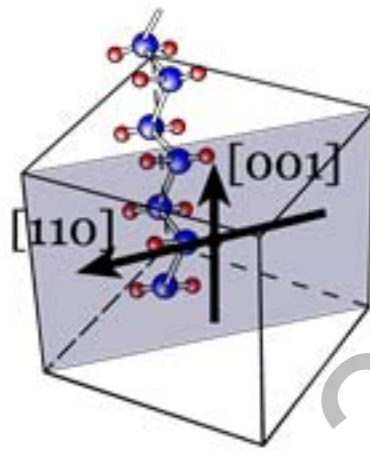
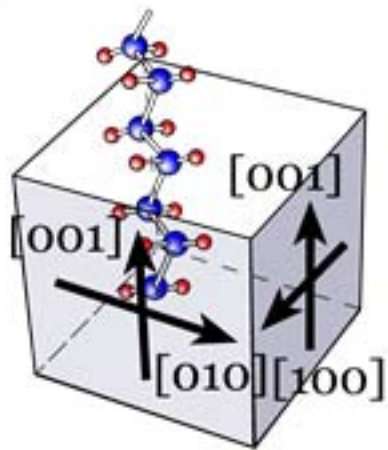
MODELLING APPROACH

Idealized 100% crystalline polyethylene



DEFORMATION MECHANISMS

Crystallographic slip systems



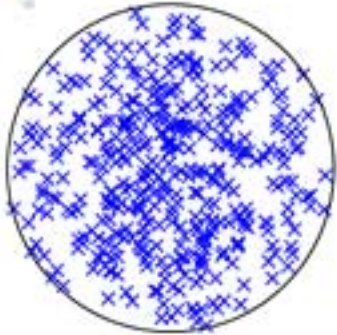
Critical shear strength	
Slip system	τ^c (MPa)
$(100) [001]$	7.2
$(010) [001]$	15.6
$\{110\} [001]$	13.0
$(100) [010]$	12.3

Bartczak *et al*, 1992

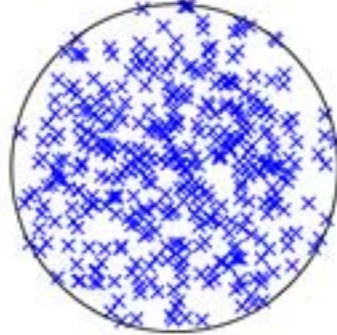
CRYSTAL AGGREGATE

Crystallographic texture

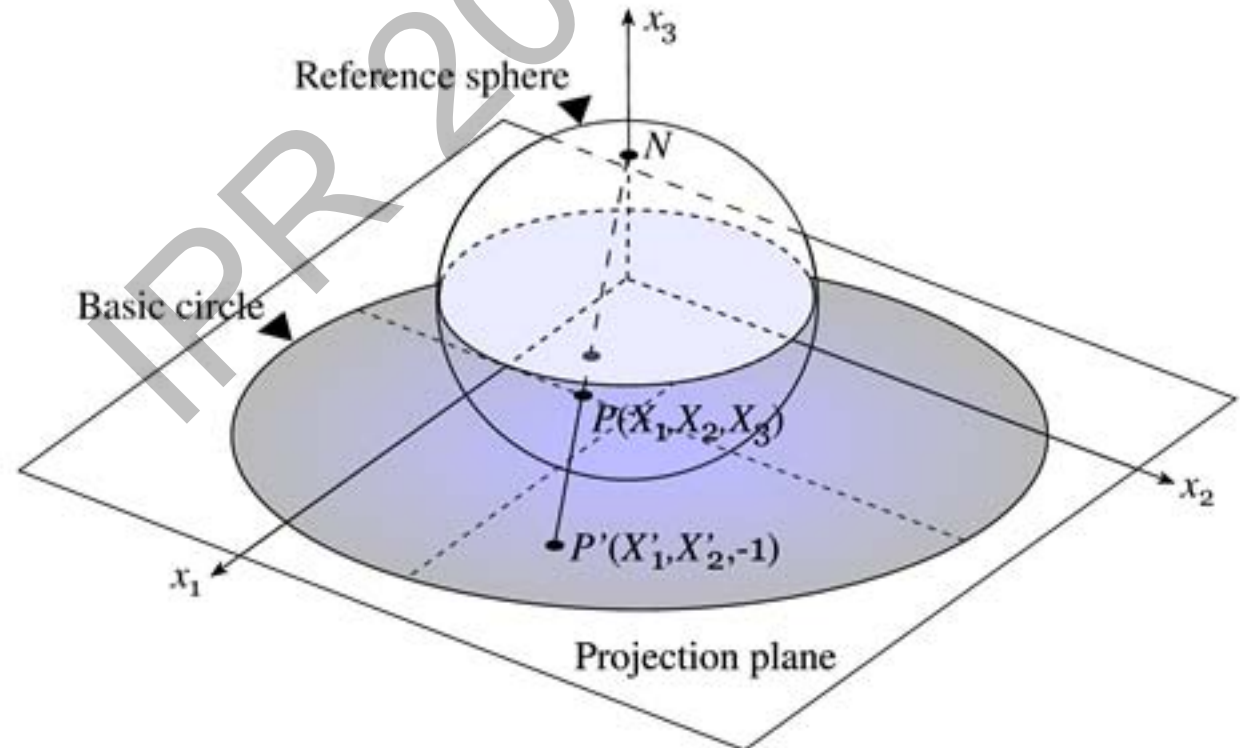
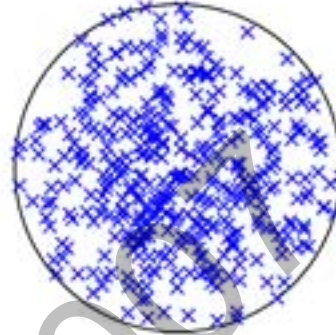
(100)



(010)



(001)



CRYSTALLINE POLYETHYLENE

Summary - Viscoplastic model

Schmid tensor

$$\mathbf{T}^\alpha = \mathbf{n}^\alpha \otimes \mathbf{s}^\alpha$$

EQUILIBRIUM
Resolved shear stress

$$\mathbf{S} = \boldsymbol{\sigma} - \text{tr}(\boldsymbol{\sigma})\mathbf{I}$$

$$\tau^\alpha = \mathbf{S} : \mathbf{T}^\alpha$$

Stress-strain relationship

$$\dot{\boldsymbol{\epsilon}}^\alpha = \dot{\boldsymbol{\epsilon}}_0^\alpha \left(\frac{\tau^\alpha}{g^\alpha} \right)^n$$

$\dot{\boldsymbol{\epsilon}}^\alpha$ Shear strain rate
 τ^α Resolved shear stress

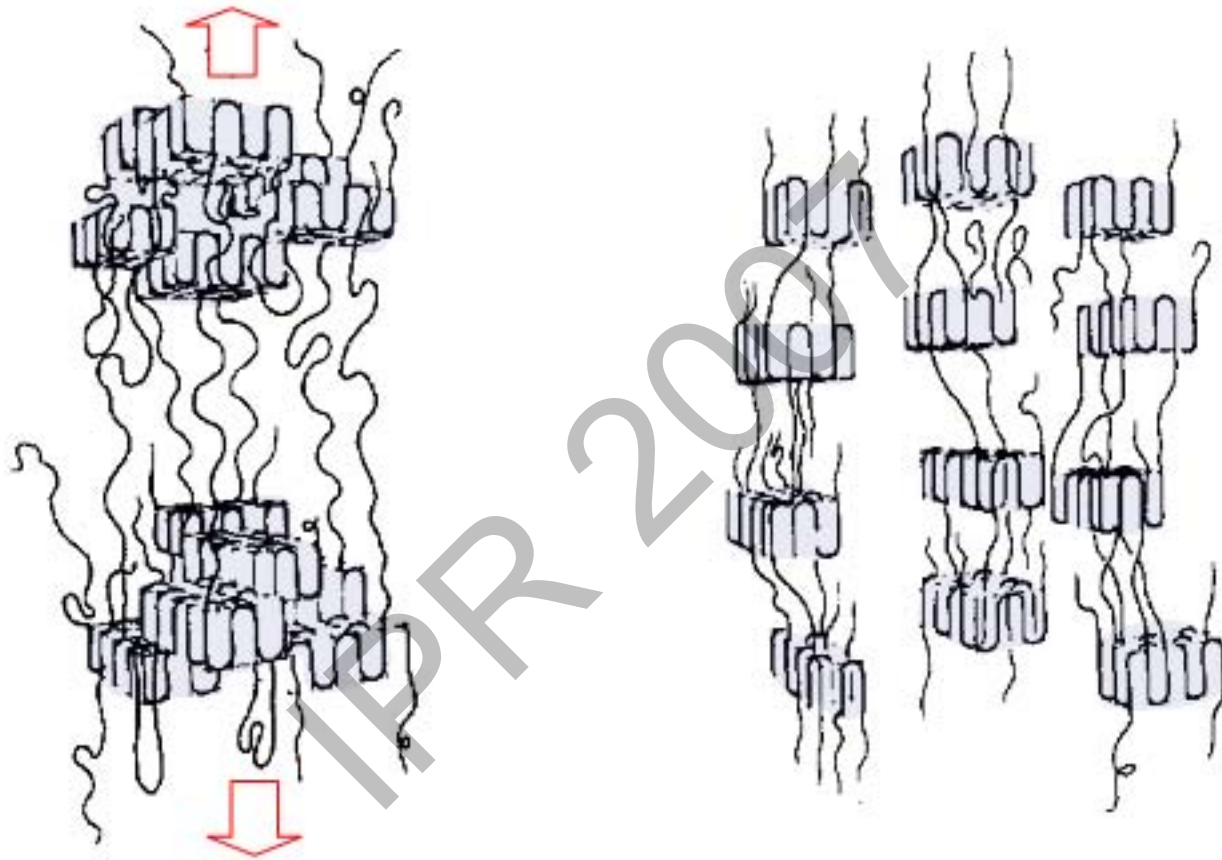
COMPATIBILITY
Deformation rate and spin

$$\mathbf{D}^p = \sum_\alpha \dot{\boldsymbol{\epsilon}}^\alpha \text{sym}(\mathbf{T}^\alpha) \quad \mathbf{W}^p = \sum_\alpha \dot{\boldsymbol{\epsilon}}^\alpha \text{skw}(\mathbf{T}^\alpha)$$

\mathbf{D}^p Deformation rate
 \mathbf{W}^p Spin

DAMAGE MODEL

Stress-strain state in single crystals

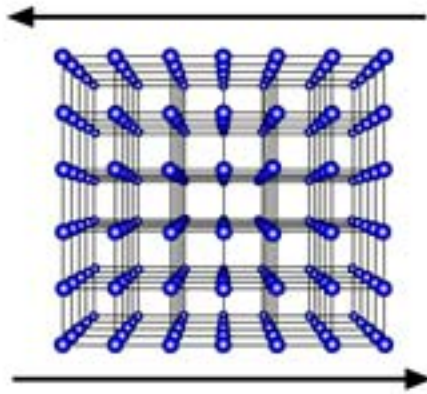


$$\dot{\epsilon}^{\alpha} = \dot{\epsilon}_0^{\alpha} \left(\frac{\tau^{\alpha}}{g^{\alpha}} \right)^n$$

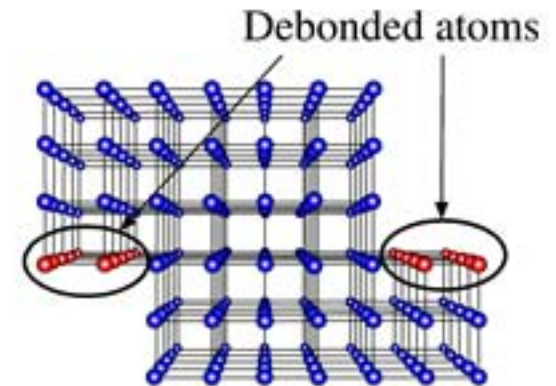
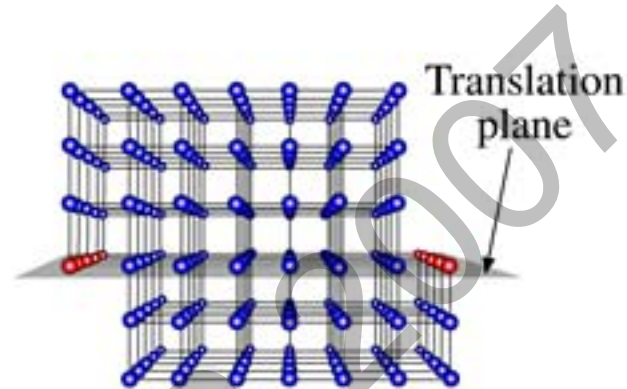
Stress-strain relationship with no damage

DAMAGE MODEL

Stress-strain state in single crystals



Undamaged crystal



Damage crystal

$$\Omega^{\alpha} = \frac{\text{Current number of atomic debonds}}{\text{Initial number of atomic bonds}} = \frac{\text{Total area of defects}}{\text{Initial area}}$$

DAMAGE MODEL

Stress-strain state in single crystals

$$\dot{\gamma}^{\alpha} = \dot{\gamma}_0^{\alpha} \left(\frac{\tau^{\alpha}}{g^{\alpha}} \right)^n$$

Classic
viscoplastic model



$$\dot{\gamma}^{\alpha} = \dot{\gamma}_0^{\alpha} \left(\frac{\tau^{\alpha}}{(1 - \Omega^{\alpha}) g^{\alpha}} \right)^n$$

Stress-strain
relationship

$$\frac{d\Omega^{\alpha}}{d\tau^{\alpha}} = \dot{\gamma}_0^{\alpha} \left| \frac{\tau^{\alpha}}{(1 - \Omega^{\alpha}) g^{\alpha}} \right|^m$$

Damage evolution
law

$$\frac{dg^{\alpha}}{d\gamma^{\alpha}} = h_0 \operatorname{sech}^2 \left(\frac{h_0}{c} \hat{\gamma}^{\alpha} \right)$$

Hardening
law

Damage model

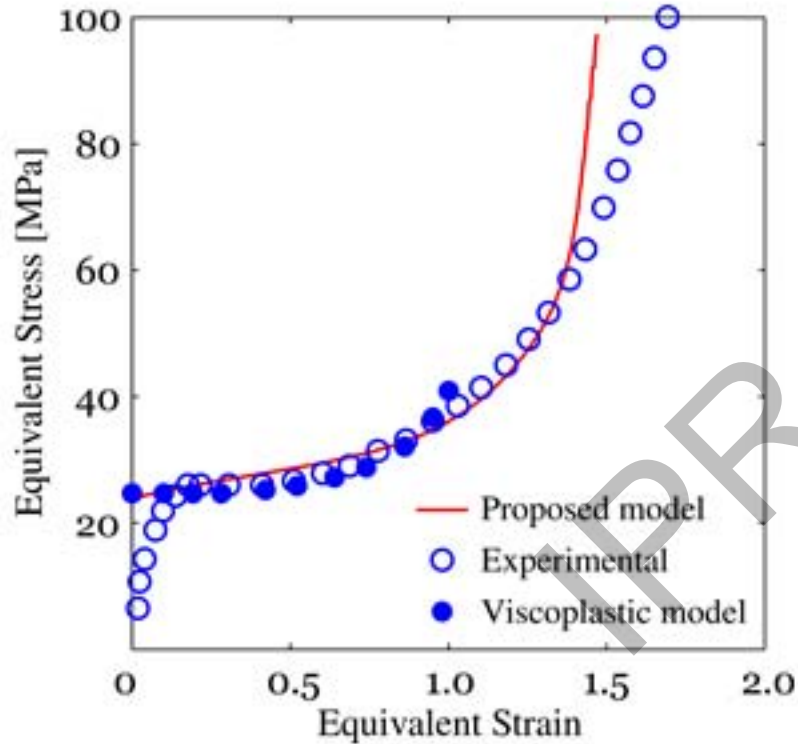
PRELIMINARY RESULTS

Proposed damage model

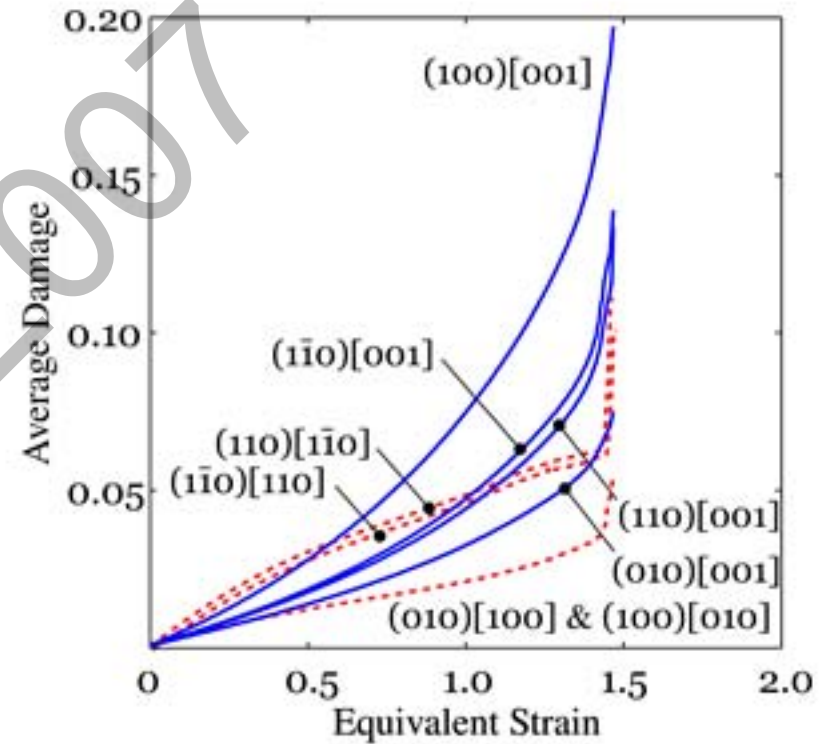
- Idealized 100% crystalline polyethylene.
- 100 randomly oriented crystals.
- Initially isotropic texture.
- Uniaxial tension and simple shear

UNIAXIAL TENSION

Idealized 100% crystalline



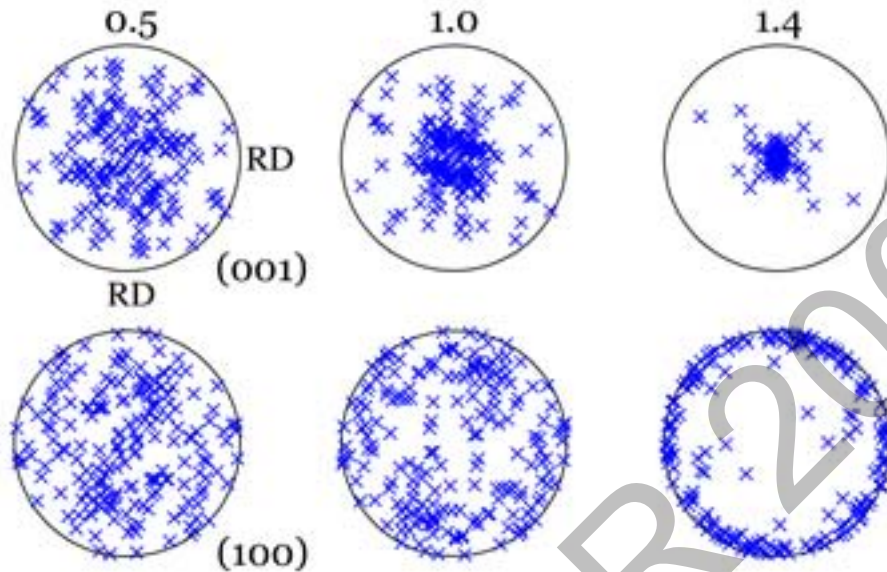
Stress-strain behaviour



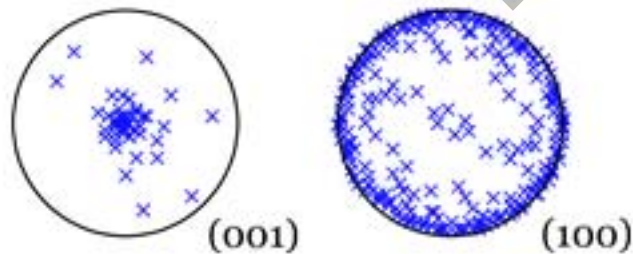
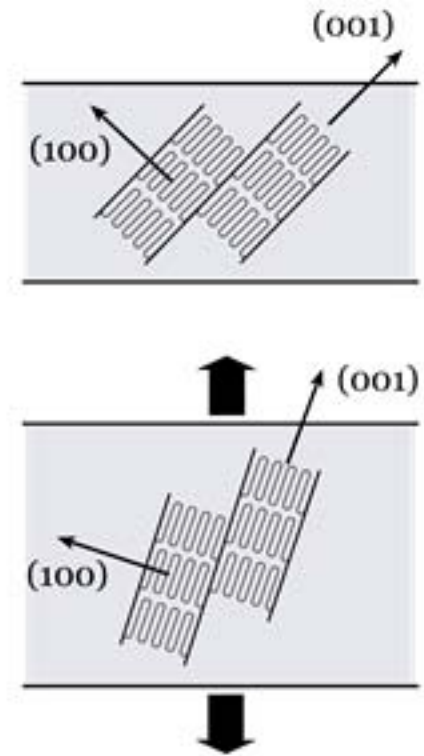
Slip systems damage

UNIAXIAL TENSION

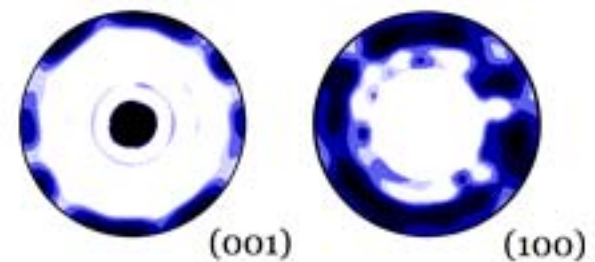
Crystallographic textures



Projection perpendicular to the loading direction



Parks & Ahzi (1.0)



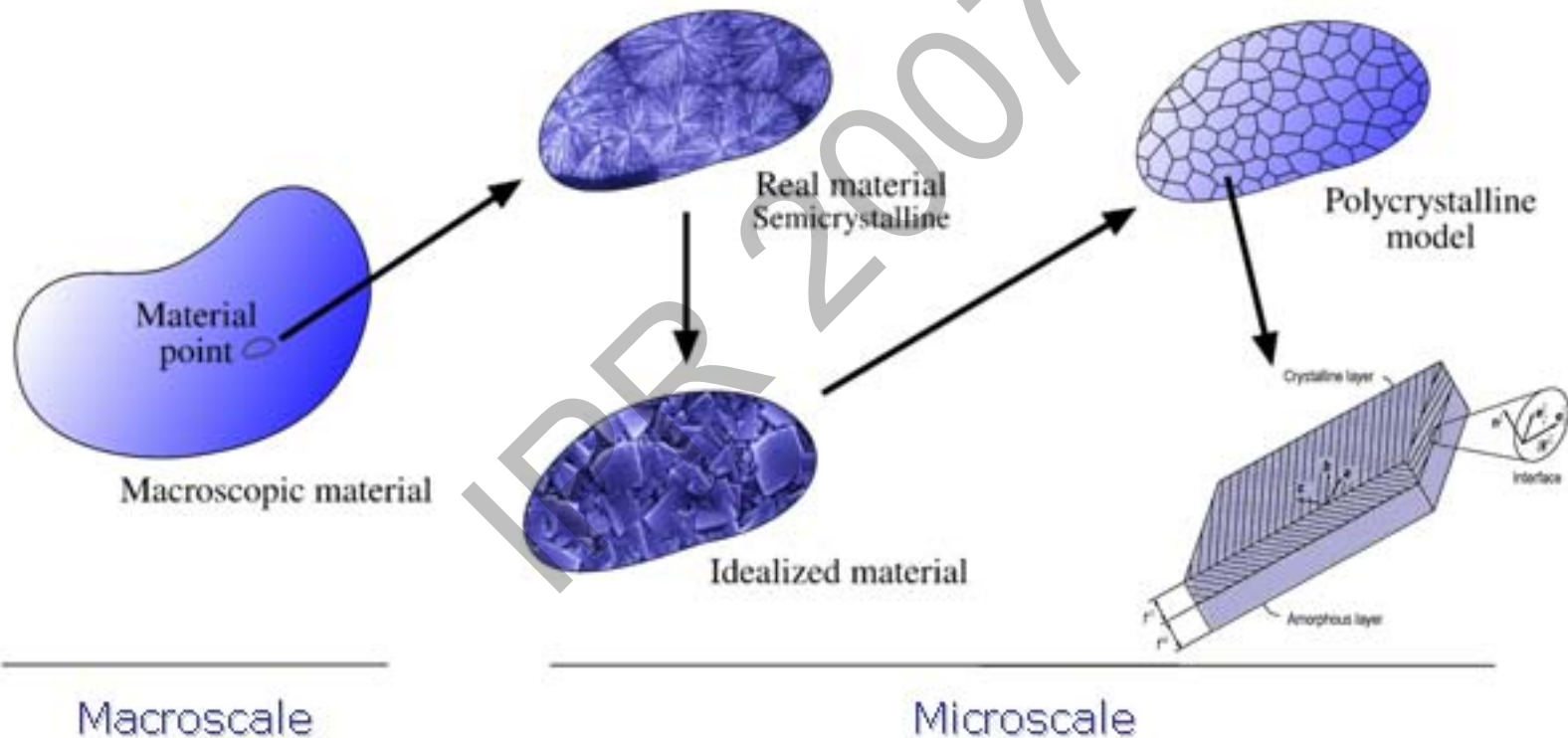
Experimental (1.8)

PRESENTATION OUTLINE

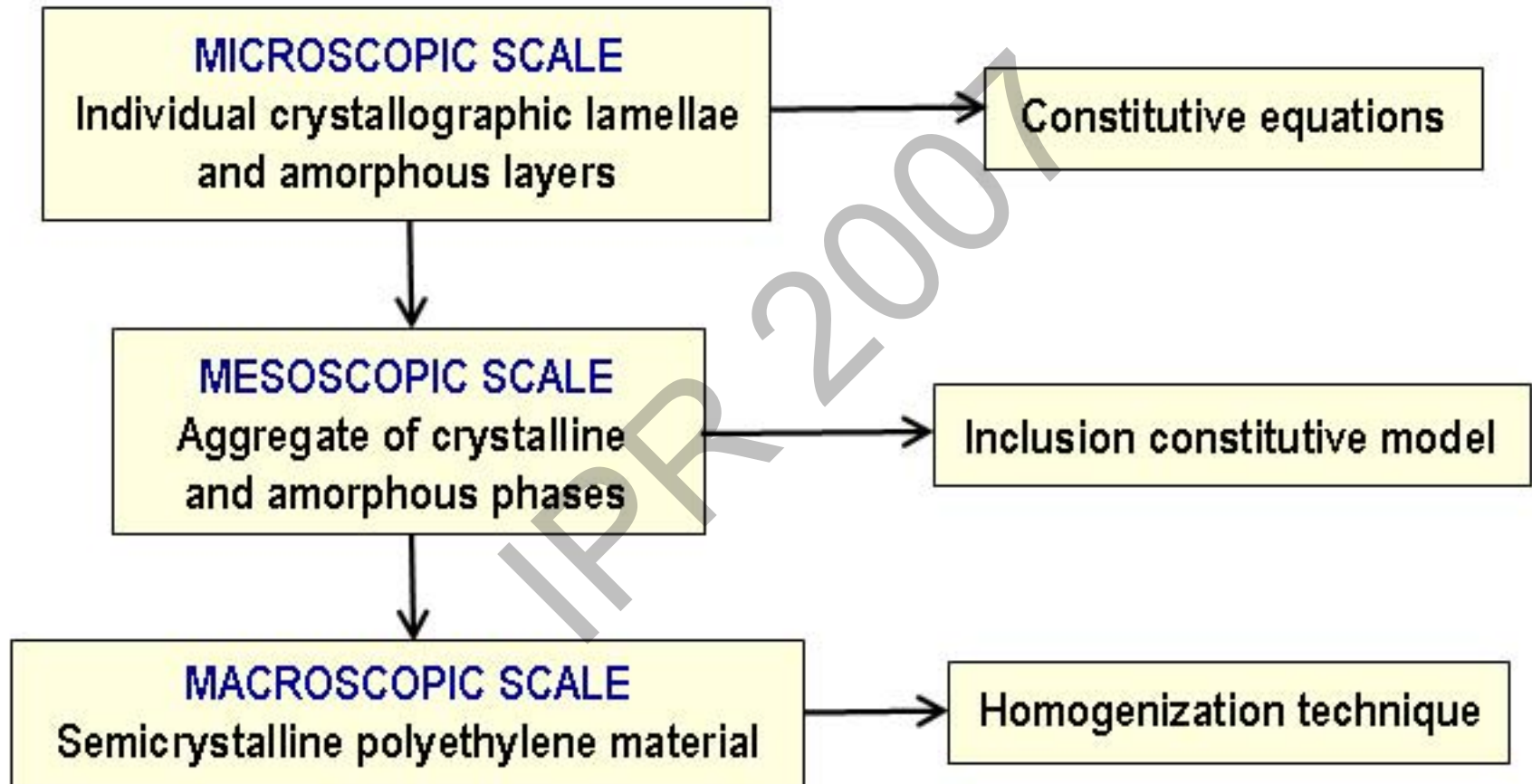
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MODELLING APPROACH

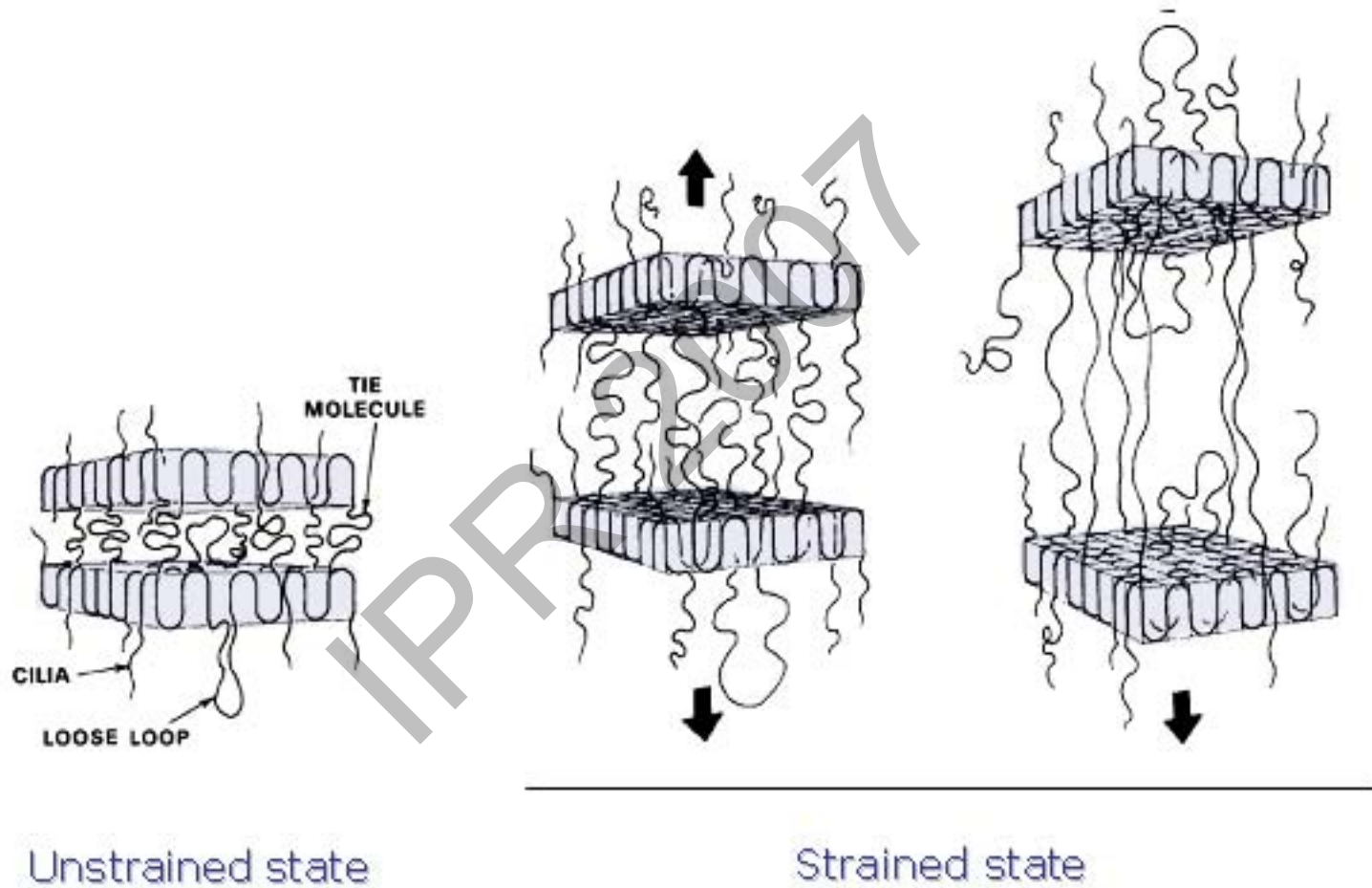
Semicrystalline polyethylene



MULTISCALE MODELLING

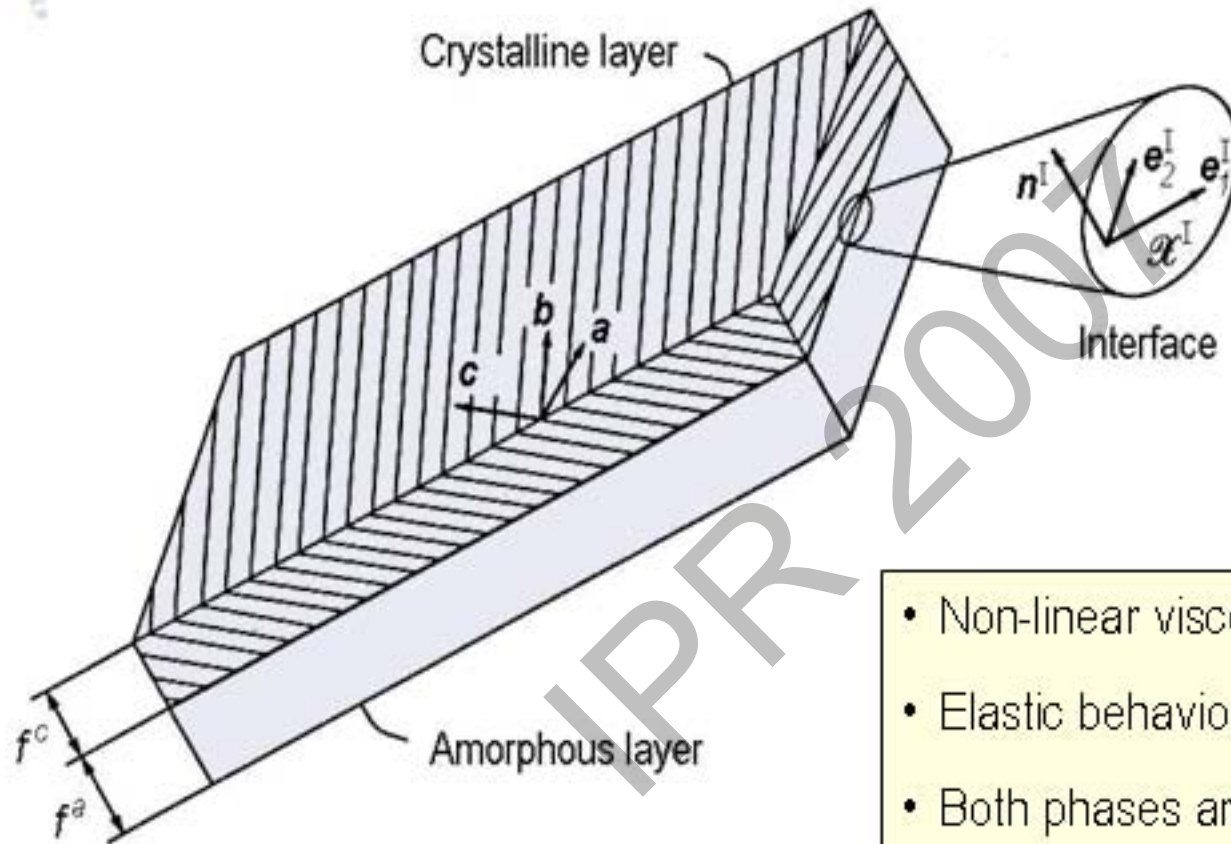


AMORPHOUS POLYETHYLENE



COMPOSITE INCLUSION

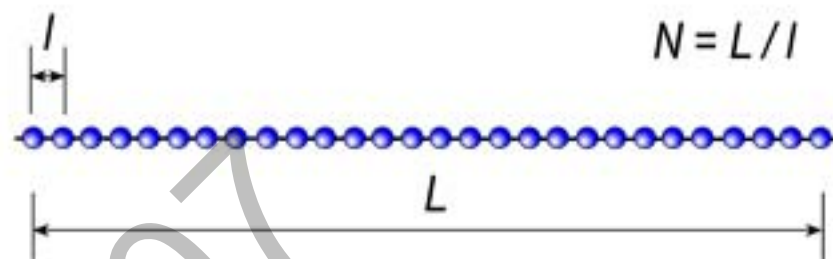
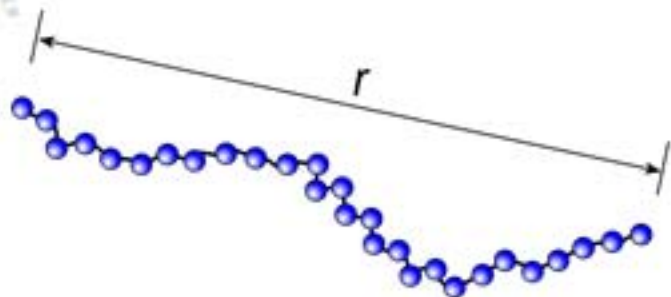
Inclusion constitutive model



- Non-linear viscoplastic behaviour.
- Elastic behaviour is not included.
- Both phases are incompressible.
- Insensitive to pressure deformation
- Deformation and stress are uniform.

AMORPHOUS POLYETHYLENE

Molecular chain entropy



Boltzmann's Principle

$$S = k \ln W$$

$$k = 1.38 \times 10^{-23} \text{ JK}^{-1}$$

Probability density of a molecular chain

$$\ln W = \text{constant} - N \left(\frac{r}{Nl} \beta + \ln \frac{\beta}{\sinh \beta} \right)$$

$$\beta = \mathcal{L}^{-1}(r/Nl)$$

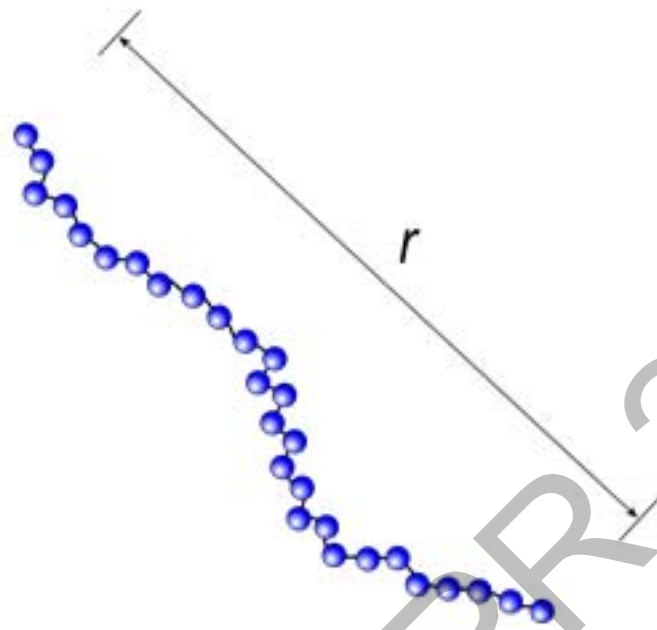
$$\mathcal{L}(\beta) = \coth \beta - 1/\beta$$

Molecular chain entropy

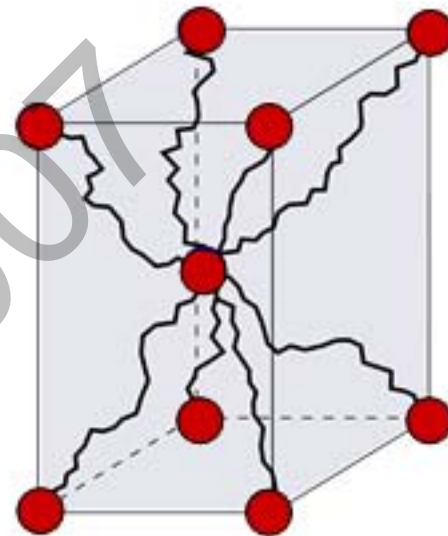
$$S = \text{constant} - kN \left(\frac{r}{Nl} \beta + \ln \frac{\beta}{\sinh \beta} \right)$$

AMORPHOUS POLYETHYLENE

Kinematic hardening – Back-stress tensor



Idealized molecule



The eight-chain model

Back-stress tensor

$$\mathbf{H} = \frac{2}{J} \mathbf{B} \cdot \frac{\partial \Psi}{\partial \mathbf{B}}$$

Helmholtz free energy $d\Psi = -T dS$

$$\mathbf{H} = \frac{nkT}{3} \sqrt{\frac{3N}{l_1}} \mathcal{L}^{-1} \left(\sqrt{\frac{l_1}{3N}} \right) \mathbf{B}$$

AMORPHOUS POLYETHYLENE

Summary – Viscoplastic model

Deviatoric Cauchy stress

$$\mathbf{S} = \boldsymbol{\sigma} - \text{tr}(\boldsymbol{\sigma})\mathbf{I}$$

Deviatoric back-stress

$$\mathbf{H}' = \frac{nkT}{3} \sqrt{\frac{3N}{I_1}} \mathcal{L}^{-1} \left(\sqrt{\frac{I_1}{3N}} \right) \left(\mathbf{B} - \frac{1}{3}\mathbf{I} \right)$$

Resolved shear stress

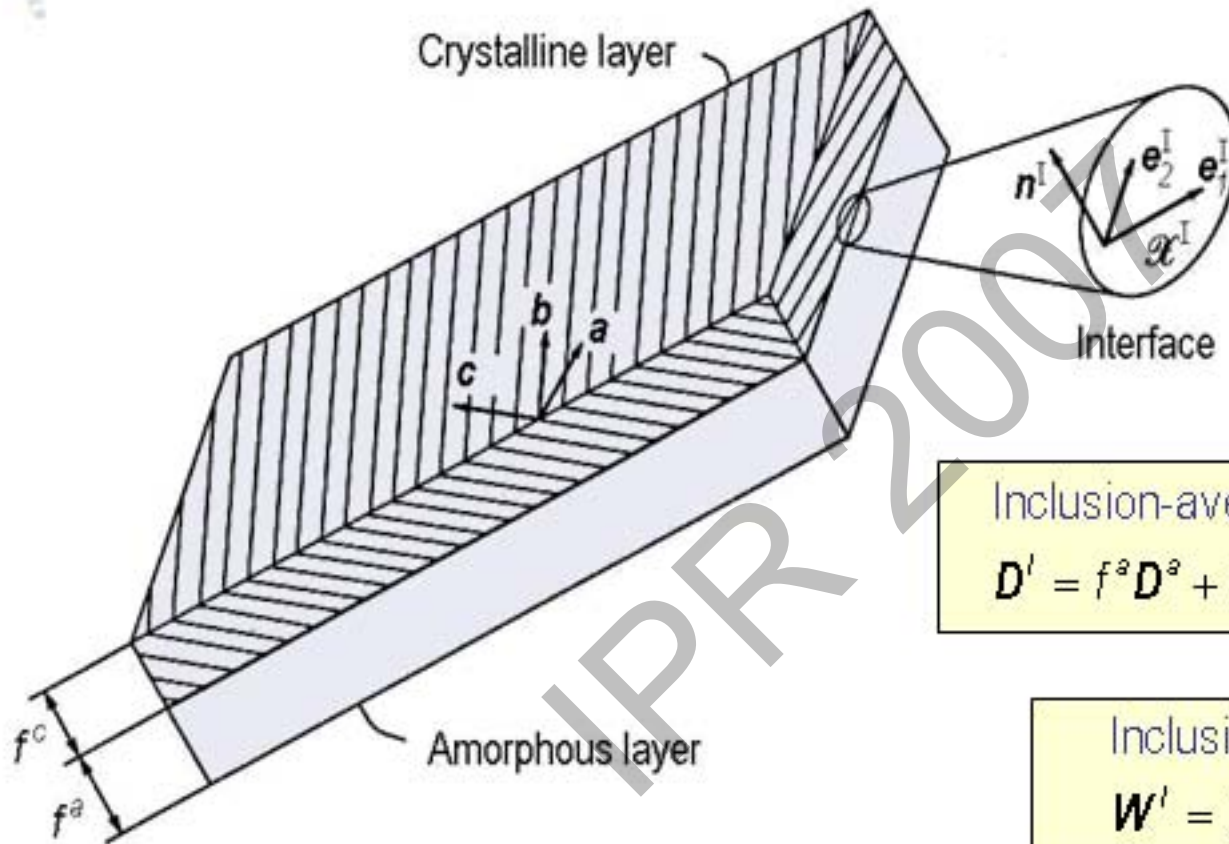
$$\tau = \left[\frac{1}{2} (\mathbf{S} - \mathbf{H}') : (\mathbf{S} - \mathbf{H}') \right]^{1/2}$$

Deformation rate

$$\mathbf{D}^p = \frac{\tau}{\tau_0} \left(\frac{\tau}{\tau_0} \right)^{m-1} \left(\frac{\mathbf{S} - \mathbf{H}'}{\tau_0} \right)$$

INCLUSION VOLUME-AVERAGING

Homogenization technique



Inclusion-averaged deformation rate

$$D^I = f^a D^a + f^c D^c = (1 - f^c) D^a + f^c D^c$$

Inclusion-averaged spin

$$W^I = (1 - f^c) W^a + f^c W^c$$

Inclusion-averaged deviatoric stress

$$S^I = (1 - f^c) S^a + f^c S^c$$

INTERACTION RELATIONS

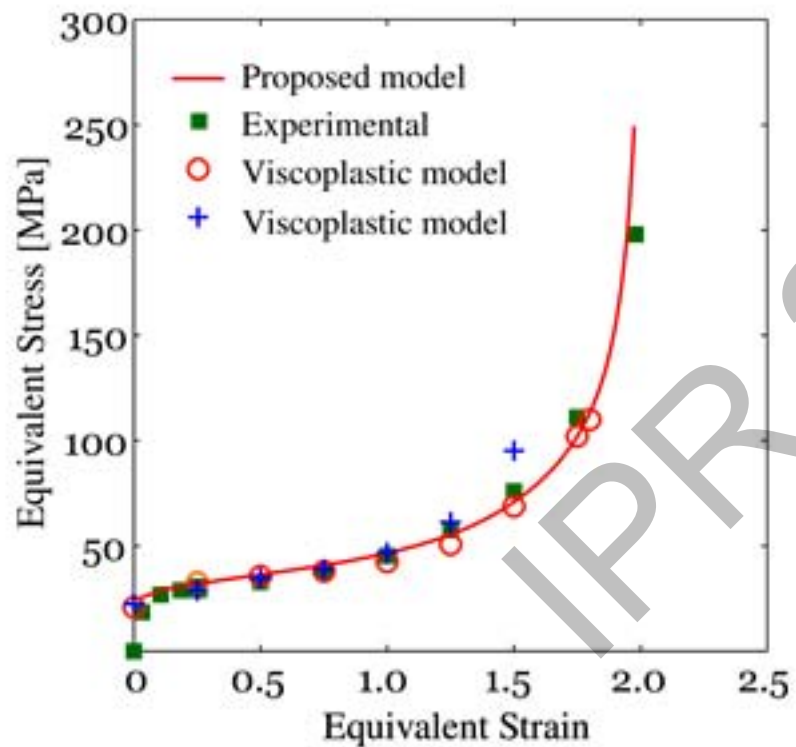
Homogenization technique

•Sachs-like Model

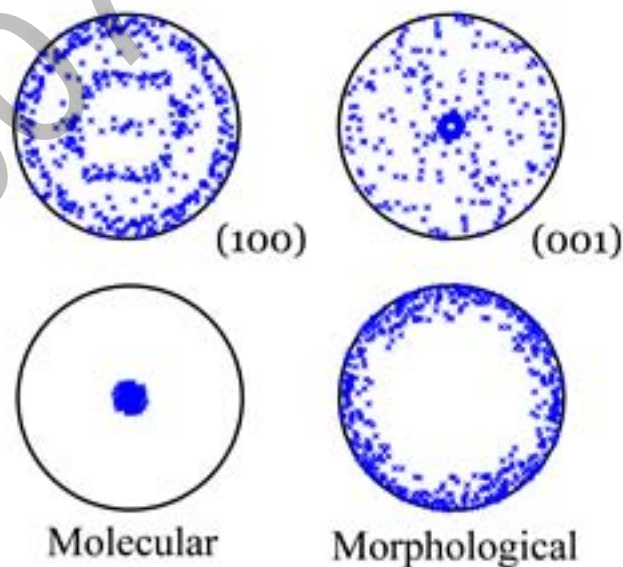
- Local and global equilibrium are satisfied by considering the stress as uniform in all component parts.
- Global compatibility is enforced as a global volume average.
- Local compatibility of the inclusion deformation with the surrounding is not addressed.

UNIAXIAL TENSION

Semi-crystalline



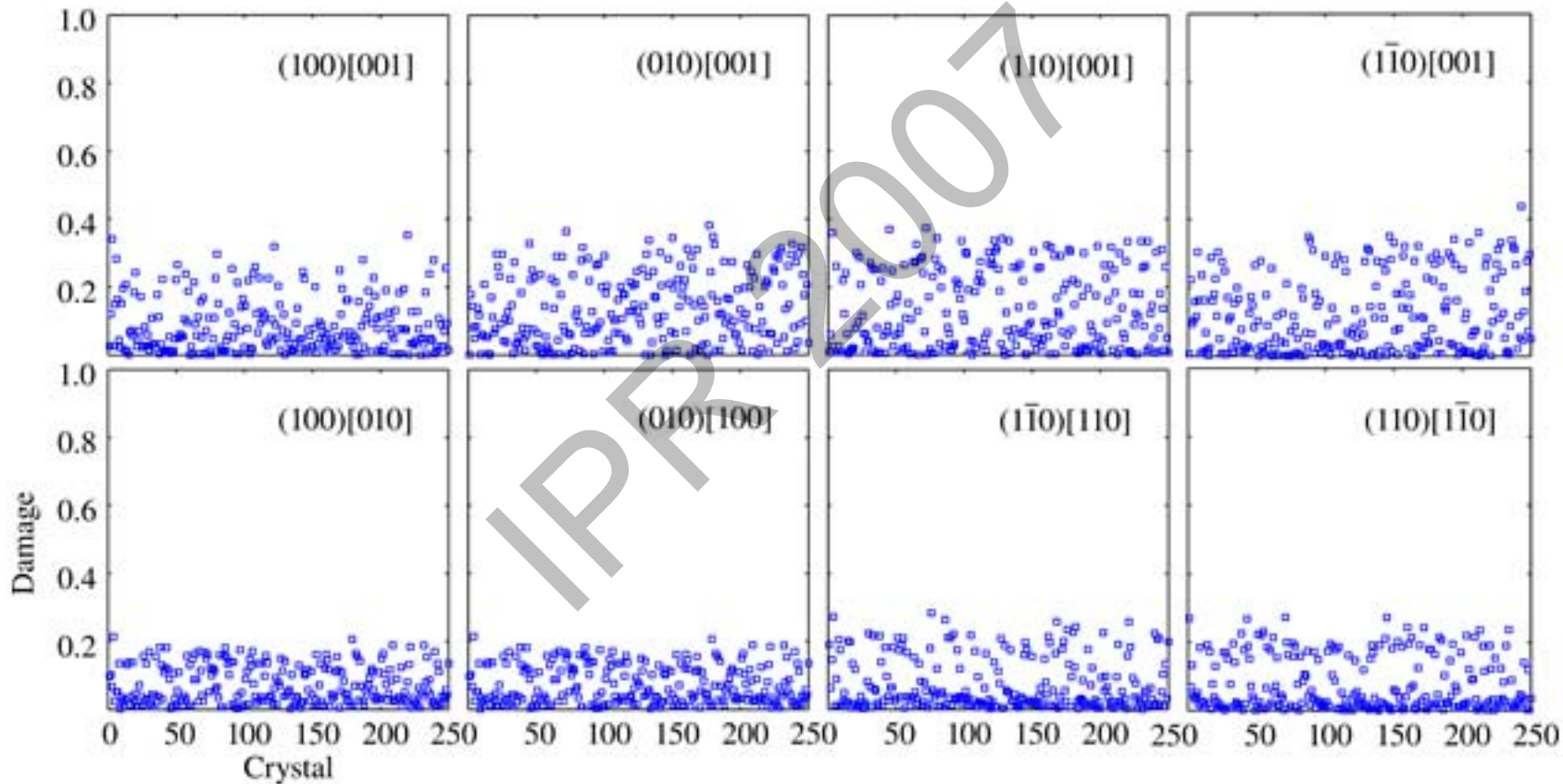
Stress-strain behaviour



Texture

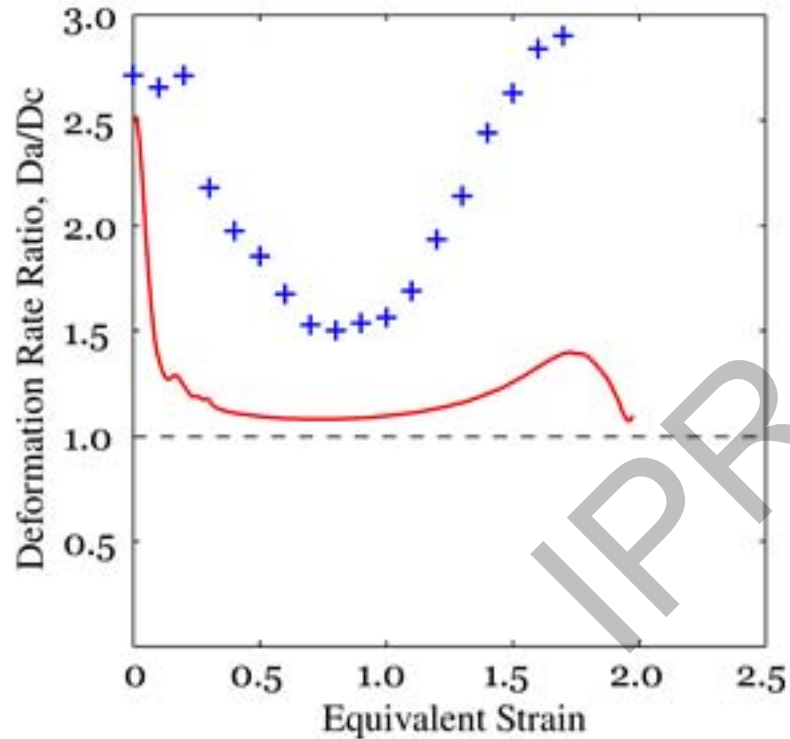
UNIAXIAL TENSION

Semi-crystalline – Slip system damage

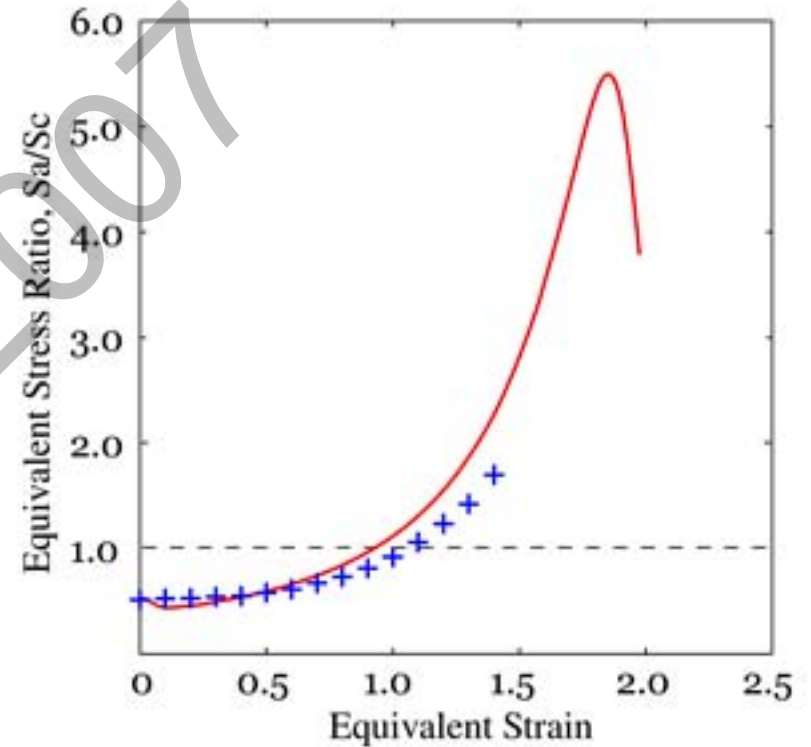


UNIAXIAL TENSION

Semi-crystalline



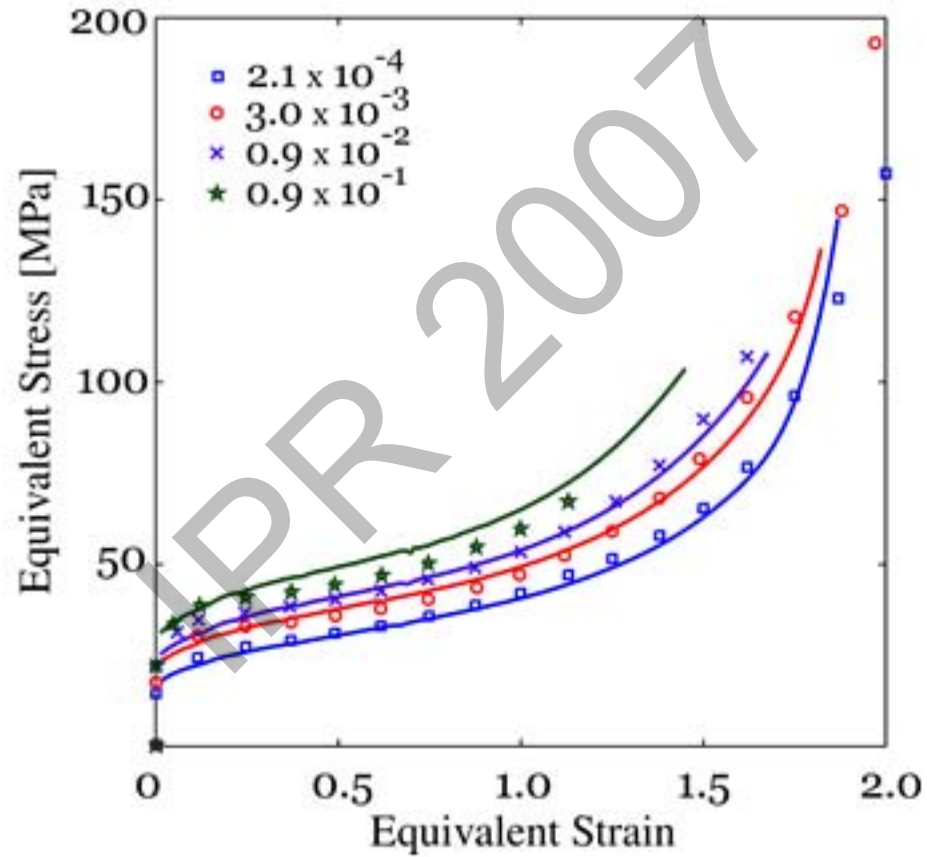
Deformation rate behaviour



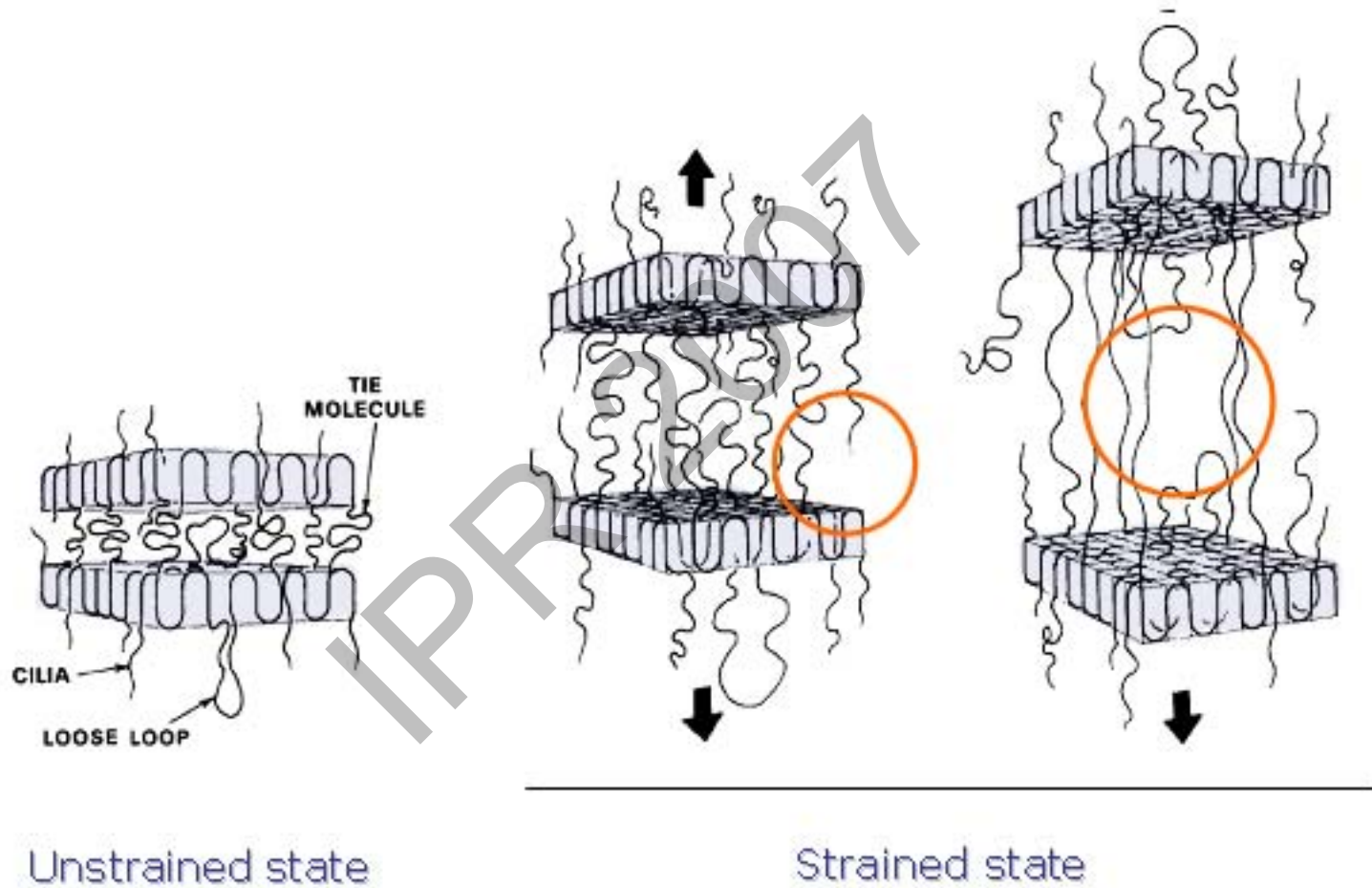
Stress behaviour

UNIAXIAL TENSION

Behaviour at different strain rates



AMORPHOUS POLYETHYLENE



AMORPHOUS POLYETHYLENE

Damage-coupled model

Deviatoric Cauchy stress

$$\mathbf{S} = \boldsymbol{\sigma} - \text{tr}(\boldsymbol{\sigma})\mathbf{I}$$

Deviatoric back-stress

$$\mathbf{H}' = \frac{nkT}{3} \sqrt{\frac{3N}{I_1}} \mathcal{L}^{-1} \left(\sqrt{\frac{I_1}{3N}} \right) (\mathbf{B} - \frac{1}{3}\mathbf{I})$$

Equivalent stress

$$\sigma = \left[\frac{3}{2} \left(\frac{\mathbf{S}}{(1-\Omega^a)} - \mathbf{H}' \right) : \left(\frac{\mathbf{S}}{(1-\Omega^a)} - \mathbf{H}' \right) \right]^{1/2}$$

Deformation rate

$$D^p = \mathfrak{E} \left(\frac{\sigma}{\sigma_0} \right)^{m-1} \left(\frac{\frac{\mathbf{S}}{(1-\Omega^a)} - \mathbf{H}'}{\sigma_0} \right)$$

Damage evolution law

$$\Omega = \Omega_{cr} \frac{W - W_0}{W} = \Omega_{cr} [1 - \exp(\Delta s/k)]$$

PRESENTATION OUTLINE

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- Modelling of semicrystalline polyethylene
- **Conclusions**

CONCLUSIONS

- **Two micromechanically based composite models to study plastic deformation and texture evolution in initially isotropic polyethylene have been proposed.**
- The models have been tested in uniaxial tension and simple shear.
- The predicted results agree with the experimental observations of macroscopic stress-strain behaviour and texture evolution in nearly all aspects.
- The models can also be applied to study plastic deformation and texture evolution in other semicrystalline polymers.

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